

ENTROPY DISSIPATION SEMI-DISCRETIZATION SCHEMES FOR FOKKER-PLANCK EQUATIONS

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ABSTRACT. We propose a new semi-discretization scheme to approximate nonlinear Fokker-Planck equations, by exploiting the gradient flow structures with respect to the optimal transport metric (2-Wasserstein metric). We discretize the underlying state by a finite graph and define a discrete 2-Wasserstein metric. Based on such metric, we introduce a dynamical system, which is a gradient flow of the discrete free energy. We prove that the new scheme maintains dissipativity of the free energy and converges to a discrete Gibbs measure at exponential (dissipation) rate. We exhibit these properties on several numerical examples.

1. INTRODUCTION

In this paper we give new semi-discretization schemes for a certain type of partial differential equations (PDEs) [15], which are gradient flows from the viewpoint of optimal transportation theory [1, 2, 3, 10, 20, 21, 24, 25].

First, let us briefly review some known facts. Consider a nonlinear Fokker-Planck equation [4, 10]

$$\frac{\partial \rho}{\partial t} = \nabla \cdot [\rho \nabla(V(x) + \int_{\mathbb{R}^d} W(x, y) \rho(t, y) dy)] + \beta \Delta \rho. \quad (1)$$

Here, the unknown $\rho(t, \cdot)$ is a probability density function supported on \mathbb{R}^d , which is thus positive and conserves the total probability. The functions $V : \mathbb{R}^d \rightarrow \mathbb{R}$, and $W : \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}$ are smooth and further $W(x, y) = W(y, x)$ for any $x, y \in \mathbb{R}^d$.

To (1) is associated a functional $\mathcal{F} : \mathcal{P}(\mathbb{R}^d) \rightarrow \mathbb{R}$, called free energy

$$\mathcal{F}(\rho) = \int_{\mathbb{R}^d} V(x) \rho(x) dx + \frac{1}{2} \int_{\mathbb{R}^d \times \mathbb{R}^d} W(x, y) \rho(x) \rho(y) dx dy + \beta \int_{\mathbb{R}^d} \rho(x) \log \rho(x) dx, \quad (2)$$

which is a summation of linear potential energy, interaction energy and linear entropy, from left to right. It is known that the free energy (2) is a Lyapunov function for (1):

$$\frac{d}{dt} \mathcal{F}(\rho(t, \cdot)) = - \int_{\mathbb{R}^d} (\nabla F(x, \rho))^2 \rho(t, x) dx \leq 0,$$

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where¹

$$F(x, \rho) := \frac{\delta}{\delta \rho(x)} \mathcal{F}(\rho) = V(x) + \int_{\mathbb{R}^d} W(x, y) \rho(y) dy + \beta \log \rho(x) + \beta.$$

Under suitable conditions on V and W , the solution $\rho(t, \cdot)$ of (1) converges to an equilibrium, called Gibbs measure².

Recent work on optimal transport treats the probability set $\mathcal{P}(\mathbb{R}^d)$ as a “manifold” equipped with the 2-Wasserstein metric. From this viewpoint, (1) is a gradient flow of the free energy $\mathcal{F}(\rho)$ on $\mathcal{P}(\mathbb{R}^d)$, see [2, 24, 25]. Furthermore, under some requirements on W , Carrillo, McCann and Villani show that $\rho(t, \cdot)$ converges to a Gibbs measure with exponential rate, see [10].

In this paper, we consider a similar matter in the discrete setting. In other words, we plan to derive a semi-discretization scheme³ for (1), which also has a gradient flow structure with respect to a discrete 2-Wasserstein metric.

Notice that the underlying state of (1) can be a variety other than \mathbb{R}^d . For instance, the domain can be a bounded open set, with a zero-flux conditions or periodic conditions. In this paper, we use a graph setting to consider all these cases.

We shall consider a graph $G = (V, E)$ to discretize the spatial domain, where V is the vertex set

$$V = \{1, 2, \dots, n\},$$

and E is the edge set. The adjacency set of the vertex $i \in V$ is denoted by

$$N(i) = \{j \in V \mid (i, j) \in E\}.$$

Here $i \in V$ represents a point in \mathbb{R}^d , and (i, j) is shorthand for an edge connecting i and j . For concreteness, we can think of G as a lattice corresponding to a uniform discretization of the domain with spacing Δx .

We consider a discrete probability set⁴ supported on all vertices:

$$\mathcal{P}(G) = \{\rho = (\rho_i)_{i=1}^n \in \mathbb{R}^n \mid \sum_{i=1}^n \rho_i = 1, \rho_i \geq 0, i \in V\}.$$

Moreover, we consider a discrete free energy of $\mathcal{F}(\rho)$, as an analog of (2)

$$\mathcal{F}(\rho) = \sum_{i=1}^n v_i \rho_i + \frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n w_{ij} \rho_i \rho_j + \beta \sum_{i=1}^n \rho_i \log \rho_i,$$

where $(v_i)_{i=1}^n = (V(i))_{i=1}^n$ and $(w_{ij})_{1 \leq i, j \leq n} = (W(i, j))_{1 \leq i, j \leq n}$ are a fixed vector, and a fixed symmetric matrix, respectively.

¹ $\frac{\delta}{\delta \rho(x)} \mathcal{F}$ is the first variation of the scalar functional \mathcal{F} .

² $\rho^*(x)$ is a Gibbs measure if it solves the fixed point problem

$$\rho^*(x) = \frac{1}{K} e^{-\frac{V(x) + \int_{\mathbb{R}^d} W(x, y) \rho^*(y) dy}{\beta}}, \quad \text{where } K = \int_{\mathbb{R}^d} e^{-\frac{V(x) + \int_{\mathbb{R}^d} W(x, y) \rho^*(y) dy}{\beta}} dx.$$

³We discretize the spatial variable, not the time variable.

⁴We abuse the notation. We let ρ_i represent a discrete probability measure, instead of probability density.

By this setting, we will derive and analyze the gradient flow of discrete free energy on $\mathcal{P}(G)$ with a “discrete” 2-Wasserstein metric. Delaying the derivation details until section 2, we show the semi-discretization directly. We propose to take

$$\frac{d\rho_i}{dt} = \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j (F_j(\rho) - F_i(\rho))_+ - \sum_{j \in N(i)} \rho_i (F_i(\rho) - F_j(\rho))_+ \right\}, \quad (3)$$

where $i \in V$, $(\cdot)_+ = \max\{\cdot, 0\}$ and

$$F_i(\rho) = \frac{\partial}{\partial \rho_i} \mathcal{F}(\rho), \quad \text{for any } i \in V,$$

and thus

$$F_i(\rho) = v_i + \sum_{j=1}^n w_{ij} \rho_j + \beta \log \rho_i + \beta.$$

Besides showing that (3) is a well defined ordinary differential equation (ODE), we demonstrate that (3) has a gradient flow structure. Firstly, the free energy is a Lyapunov function of (3):

$$\frac{d}{dt} \mathcal{F}(\rho(t)) = - \sum_{(i,j) \in E} \left(\frac{F_i(\rho) - F_j(\rho)}{\Delta x} \right)_+^2 \rho_i \leq 0.$$

Then, if $\rho(t)$ converges to an equilibrium ρ^∞ , then we will show that such equilibrium is a discrete Gibbs measure

$$\rho_i^\infty = \frac{1}{K} e^{-\frac{v_i + \sum_{j=1}^n w_{ij} \rho_j^\infty}{\beta}}, \quad K = \sum_{i=1}^n e^{-\frac{v_i + \sum_{j=1}^n w_{ij} \rho_j^\infty}{\beta}}.$$

Furthermore, if ρ^∞ is a strictly local minimizer of the free energy, and $\rho(t)$ is in its basin of attraction for the gradient dynamics, then we will show that the convergence speed is exponential:

$$\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty) \leq e^{-Ct} (\mathcal{F}(\rho^0) - \mathcal{F}(\rho^\infty)),$$

where C is a positive constant. In fact, we will say more about this convergence. We will give an explicit formula for the asymptotic convergence rate, which mimics the role of the Hessian of the free energy at the Gibbs measure.

Finally, we will show that (3) is a consistent scheme for PDE (1). In fact, we will derive a general consistent scheme for general diffusion systems, which may not be gradient flows.

The semi-discretization scheme in this paper is largely inspired by [12, 16], the upwind scheme of [5], and optimal transport theory [25]. In addition, the convergence result is influenced by the work of Carrillo, McCann and Villani, [10]. Our method can be viewed as a discrete entropy⁵ dissipation method [11, 18], with a dynamical twist.

In the literature, people have studied 2-Wasserstein metric and Fokker-Planck equations in discrete settings for a long time [7, 17, 14, 19, 22, 8, 9]. Maas [17] and Mielke [19] introduce a different discrete 2-Wasserstein metric. Based on such metric, they analyze the convergence rate of some schemes for one-dimensional linear Fokker-Planck equations. Our scheme shows exponential convergence for all linear and nonlinear cases. Carrillo,

⁵Here the entropy refers the difference of free energy

$$\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty).$$

Chertock, Huang, Wolansky [8, 9] have recently designed several algorithms based on optimal transport viewpoint. Particularly, the first order scheme designed in [8] share some similarities with (3) for a lattice graph. However, (3) can be used in a broader context (e.g., a manifold approximated by a finite graph). In addition, our derivation is from the viewpoint of the discrete 2-Wasserstein metric.

This paper is arranged as follows. In section 2, we derive (3) based on a discrete 2-Wasserstein metric. With respect to this metric, (3)'s gradient flow properties are given. In section 3, we show that the solution of (3) converges to a discrete Gibbs measure exponentially fast. Numerical analysis and several experiments on (3) are discussed in sections 4 and 5. Finally, in the Appendix, we give a few explicit models of (simple) discretizations to elucidate how the speed of convergence to a stationary measure depends on the boundary conditions.

2. SEMI-DISCRETIZATION SCHEME

In this section, we show that (3) is a gradient flow for the discrete free energy on the probability set $\mathcal{P}(G)$ in the following three steps. First, we define a discrete 2-Wasserstein metric. Second, based on such metric, we derive (3) as a gradient flow of the discrete free energy. Third, we show that (3) is a well defined ODE with gradient flow structures.

2.1. Discrete 2-Wasserstein metric.

2.1.1. *Motivation.* Recall that the 2-Wasserstein metric (Benamou-Brenier formula, [3]) is a metric defined on a probability set supported on \mathbb{R}^d :

$$W_2(\rho^0, \rho^1)^2 = \inf_{\Phi} \left\{ \int_0^1 (\nabla \Phi, \nabla \Phi)_\rho dt : \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \nabla \Phi) = 0, \rho(0) = \rho^0, \rho(1) = \rho^1 \right\},$$

where $(\cdot, \cdot)_\rho$ represents an inner product on the probability set:

$$(\nabla \Phi, \nabla \Phi)_\rho = \int_{\mathbb{R}^d} (\nabla \Phi(t, x))^2 \rho(t, x) dx,$$

and the infimum is taken among the potential functions $\Phi(t, x) \in \mathbb{R}^d$.

We give a similar metric definition on a discrete setting, which is a finite graph $G = (V, E)$. Consider a probability set supported on V with all positive measures:

$$\mathcal{P}_o(G) = \left\{ \rho = (\rho_i)_{i=1}^n \mid \sum_{i=1}^n \rho_i = 1, \rho_i > 0, \text{ for any } i \in V \right\}.$$

We use three steps to define the metric on $\mathcal{P}_o(G)$. Firstly, we define a potential vector field on graph

$$\nabla_G \Phi := (\Phi_i - \Phi_j)_{(i,j) \in E},$$

with the potential function $\Phi := (\Phi_i)_{i=1}^n$. Secondly, we introduce the divergence operator⁶ at $\rho \in \mathcal{P}_o(G)$:

$$\text{div}_G(\rho \nabla_G \Phi) := \left(-\frac{1}{\Delta x^2} \sum_{j \in N(i)} (\Phi_i - \Phi_j) g_{ij}(\rho) \right)_{i=1}^n,$$

⁶Compare with $\nabla \cdot (\rho \nabla \Phi)$.

where

$$g_{ij}(\rho) := \begin{cases} \rho_i & \text{if } F_i(\rho) > F_j(\rho), \ j \in N(i), \\ \rho_j & \text{if } F_i(\rho) < F_j(\rho), \ j \in N(i), \\ \frac{\rho_i + \rho_j}{2} & \text{if } F_i(\rho) = F_j(\rho), \ j \in N(i), \end{cases} \quad (4)$$

and $F_i(\rho) := \frac{\partial}{\partial \rho_i} \mathcal{F}(\rho)$. Lastly, we construct an inner product on $\mathcal{P}_o(G)$:

$$(\nabla_G \Phi, \nabla_G \Phi)_\rho := \frac{1}{2\Delta x^2} \sum_{(i,j) \in E} (\Phi_i - \Phi_j)^2 g_{ij}(\rho),$$

where $\frac{1}{2}$ is due to the fact that every edge in G is counted twice, i.e. $(i,j), (j,i) \in E$.

2.1.2. Discrete metric. We are ready to introduce a discrete 2-Wasserstein metric on $\mathcal{P}_o(G)$.

Definition 1. For any $\rho^0, \rho^1 \in \mathcal{P}_o(G)$, define

$$(W_2(\rho^0, \rho^1))^2 := \inf_{\Phi} \left\{ \int_0^1 (\nabla_G \Phi, \nabla_G \Phi)_\rho dt : \frac{d\rho}{dt} + \text{div}_G(\rho \nabla_G \Phi) = 0, \quad \rho(0) = \rho^0, \quad \rho(1) = \rho^1 \right\},$$

where the infimum is taken over all Φ for which ρ is a continuously differentiable curve $\rho : [0, 1] \rightarrow \mathcal{P}_o(G)$.

We justify that W_2 in Definition 1 is a well defined metric⁷. We endow $\mathcal{P}_o(G)$ with an inner product on its tangent space

$$T_\rho \mathcal{P}_o(G) = \{(\sigma_i)_{i=1}^n \in \mathbb{R}^n \mid \sum_{i=1}^n \sigma_i = 0\}.$$

Consider the equivalence relation “ \sim ” in \mathbb{R}^n which stands for “modulo additive constants,” so that the quotient space is

$$\mathbb{R}^n / \sim = \{[\Phi] \mid (\Phi_i)_{i=1}^n \in \mathbb{R}^n\}, \quad \text{where} \quad [\Phi] = \{(\Phi_1 + c, \dots, \Phi_n + c) \mid c \in \mathbb{R}^1\}.$$

We introduce an identification map

$$\tau : \mathbb{R}^n / \sim \rightarrow T_\rho \mathcal{P}_o(G), \quad \tau([\Phi]) := \left(\sum_{j \in N(i)} \frac{1}{\Delta x^2} (\Phi_i - \Phi_j) g_{ij}(\rho) \right)_{i=1}^n.$$

Lemma 2. The map $\tau : \mathbb{R}^n / \sim \rightarrow T_\rho \mathcal{P}_o(G)$ is a well defined map, linear, and one to one.

Proof. First, we show that τ is well defined. We denote

$$\sigma_i = \frac{1}{\Delta x^2} \sum_{j \in N(i)} (\Phi_i - \Phi_j) g_{ij}(\rho).$$

⁷Our approach is a discrete version of Otto calculus [25].

Our task is equivalent to show $\sum_{i=1}^n \sigma_i = 0$. Indeed,

$$\begin{aligned} \sum_{i=1}^n \sigma_i &= \frac{1}{\Delta x^2} \left\{ \sum_{i=1}^n \sum_{j \in N(i)} (\Phi_i - \Phi_j) g_{ij}(\rho) \right\} \\ &= \frac{1}{\Delta x^2} \left\{ \sum_{(i,j) \in E} \Phi_i g_{ij}(\rho) - \sum_{(i,j) \in E} \Phi_j g_{ij}(\rho) \right\} \\ &\quad \text{Relabel } i \text{ and } j \text{ on the first formula} \\ &= \frac{1}{\Delta x^2} \left\{ \sum_{(i,j) \in E} \Phi_j g_{ji}(\rho) - \sum_{(i,j) \in E} \Phi_j g_{ij}(\rho) \right\} = 0. \end{aligned}$$

Hence, the map τ is a well-defined linear map.

Next, we show τ is one to one. Since $T_\rho \mathcal{P}_o(G)$ and \mathbb{R}^n / \sim are $(n-1)$ dimensional linear spaces, we only need to prove τ is injective. I.e., if

$$\sigma_i = \frac{1}{\Delta x^2} \sum_{j \in N(i)} g_{ij}(\rho) (\Phi_i - \Phi_j) = 0, \quad \text{for any } i \in V,$$

then $[\Phi] = 0$, meaning that $\Phi_1 = \Phi_2 = \dots = \Phi_n$.

Assume this is not true. Let $c = \max_{i \in V} \Phi_i$. Since the graph G is connected, there exists an edge $(k, l) \in E$, such that $\Phi_l = c$ and $\Phi_k < c$. But, since $\sigma_l = 0$, we know that

$$\Phi_l = \frac{\sum_{j \in N(l)} g_{lj}(\rho) \Phi_j}{\sum_{j \in N(l)} g_{lj}(\rho)} = c + \frac{\sum_{j \in N(l)} g_{lj}(\rho) (\Phi_j - c)}{\sum_{j \in N(l)} g_{lj}(\rho)} < c,$$

which contradicts $\Phi_l = c$. \square

This identification map induces a scalar inner product on $\mathcal{P}_o(G)$.

Definition 3. For any two tangent vectors $\sigma^1, \sigma^2 \in T_\rho \mathcal{P}_o(G)$, we define an inner product $g : T_\rho \mathcal{P}_o(G) \times T_\rho \mathcal{P}_o(G) \rightarrow \mathbb{R}$:

$$g(\sigma^1, \sigma^2) := \frac{1}{2\Delta x^2} \sum_{(i,j) \in E} g_{ij}(\rho) (\Phi_i^1 - \Phi_j^1)(\Phi_i^2 - \Phi_j^2), \quad (5)$$

where $[\Phi^1], [\Phi^2] \in \mathbb{R}^n / \sim$, are such that $\sigma^1 = \tau([\Phi^1])$, $\sigma^2 = \tau([\Phi^2])$.

Under the above setting, we have

$$(W_2(\rho^0, \rho^1))^2 = \inf \left\{ \int_0^1 g(\sigma, \sigma) dt : \frac{d\rho}{dt} = \sigma, \quad \rho(0) = \rho^0, \quad \rho(1) = \rho^1, \quad \rho \in \mathcal{C} \right\},$$

where \mathcal{C} is the set of all continuously differentiable curves $\rho : [0, 1] \rightarrow \mathcal{P}_o(G)$. So, the metric is well defined (see more details in [12]).

2.2. Derivation of the scheme. Based on the metric manifold $(\mathcal{P}_o(G), W_2)$, we now derive the semi-discretization scheme (3) as gradient flow of the discrete free energy.

In abstract form, the gradient flow is defined by

$$\frac{d\rho}{dt} = -\text{grad}_{\mathcal{P}_o(G)} \mathcal{F}(\rho).$$

Here $\text{grad}\mathcal{F}(\rho)$ is in the tangent space $T_\rho\mathcal{P}_o(G)$, which is defined by the duality condition:

$$g(\text{grad}_{\mathcal{P}_o(G)}\mathcal{F}(\rho), \sigma) = d\mathcal{F}(\rho) \cdot \sigma, \quad \text{for any } \sigma \in T_\rho\mathcal{P}_o(G),$$

where $d\mathcal{F} \cdot \sigma = \sum_{i=1}^n \frac{\partial}{\partial \rho_i} \mathcal{F}(\rho) \sigma_i$. Hence the gradient flow satisfies

$$\left(\frac{d\rho}{dt}, \sigma \right)_\rho + d\mathcal{F}(\rho) \cdot \sigma = 0, \quad \text{for any } \sigma \in T_\rho\mathcal{P}_o(G). \quad (6)$$

Following (6), we derive (3) in Theorem 4 below.

Theorem 4. *Given a graph G , a constant $\beta > 0$, a vector $(v_i)_{i=1}^n$ and a symmetric matrix $(w_{ij})_{1 \leq i, j \leq n}$. Then the gradient flow of the discrete free energy*

$$\mathcal{F}(\rho) = \sum_{i=1}^n v_i \rho_i + \frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n w_{ij} \rho_i \rho_j + \beta \sum_{i=1}^n \rho_i \log \rho_i,$$

on the metric manifold $(\mathcal{P}_o(G), W_2)$, is

$$\frac{d\rho_i}{dt} = \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j (F_j(\rho) - F_i(\rho))_+ - \sum_{j \in N(i)} \rho_i (F_i(\rho) - F_j(\rho))_+ \right\},$$

for any $i \in V$. Here $F_i(\rho) = v_i + \sum_{j=1}^n w_{ij} \rho_j + \beta \log \rho_i + \beta$.

Proof of Theorem 4. We show the derivation of (3). For any $\sigma \in T_\rho\mathcal{P}_o(G)$, there exists $[\Phi] \in \mathbb{R}^n / \sim$, such that $\tau([\Phi]) = \sigma$. On one hand, we denote $\frac{d\rho}{dt} = (\frac{d\rho_i}{dt})_{i=1}^n$. From definition 3,

$$\begin{aligned} \left(\frac{d\rho}{dt}, \sigma \right)_\rho &= \frac{1}{2\Delta x^2} \sum_{(i,j) \in E} g_{ij}(\rho) \left(\frac{d\rho_i}{dt} - \frac{d\rho_j}{dt} \right) (\Phi_i - \Phi_j) \\ &= \frac{1}{2\Delta x^2} \left\{ \sum_{(i,j) \in E} g_{ij}(\rho) \frac{d\rho_i}{dt} (\Phi_i - \Phi_j) - \sum_{(i,j) \in E} g_{ij}(\rho) \frac{d\rho_j}{dt} (\Phi_i - \Phi_j) \right\} \\ &\quad \text{Relabel } i, j \text{ on the second formula} \\ &= \frac{1}{2\Delta x^2} \left\{ \sum_{(i,j) \in E} g_{ij}(\rho) \frac{d\rho_i}{dt} (\Phi_i - \Phi_j) - \sum_{(j,i) \in E} g_{ji}(\rho) \frac{d\rho_i}{dt} (\Phi_j - \Phi_i) \right\} \\ &= \frac{1}{\Delta x^2} \sum_{(i,j) \in E} g_{ij}(\rho) \frac{d\rho_i}{dt} (\Phi_i - \Phi_j) \\ &= \frac{1}{\Delta x^2} \sum_{i=1}^n \frac{d\rho_i}{dt} \sum_{j \in N(i)} g_{ij}(\rho) (\Phi_i - \Phi_j). \end{aligned} \quad (7)$$

At the same time, we also have

$$\begin{aligned}
d\mathcal{F}(\rho) \cdot \sigma &= \sum_{i=1}^n \frac{\partial}{\partial \rho_i} \mathcal{F}(\rho) \cdot \sigma_i = \frac{1}{\Delta x^2} \sum_{i=1}^n F_i(\rho) \sum_{j \in N(i)} g_{ij}(\rho) (\Phi_i - \Phi_j) \\
&= \frac{1}{\Delta x^2} \left\{ \sum_{i=1}^n \sum_{j \in N(i)} g_{ij}(\rho) F_i(\rho) \Phi_i - \sum_{i=1}^n \sum_{j \in N(i)} g_{ij}(\rho) F_i(\rho) \Phi_j \right\} \\
&\quad \text{Relabel } i \text{ and } j \text{ on second formula} \tag{8} \\
&= \frac{1}{\Delta x^2} \left\{ \sum_{(i,j) \in E} g_{ij}(\rho) F_i(\rho) \Phi_i - \sum_{(i,j) \in E} g_{ji}(\rho) F_j(\rho) \Phi_i \right\} \\
&= \frac{1}{\Delta x^2} \left\{ \sum_{i=1}^n \sum_{j \in N(i)} g_{ij}(\rho) (F_i(\rho) - F_j(\rho)) \Phi_i \right\}.
\end{aligned}$$

Combining (7) and (8) into (6), we have

$$\begin{aligned}
0 &= \left(\frac{d\rho}{dt}, \sigma \right)_\rho + d\mathcal{F}(\rho) \cdot \sigma \\
&= \sum_{i=1}^n \left\{ \frac{d\rho_i}{dt} + \frac{1}{\Delta x^2} \sum_{j \in N(i)} g_{ij}(\rho) (F_i(\rho) - F_j(\rho)) \right\} \Phi_i.
\end{aligned}$$

Since the above formula is true for all $(\Phi_i)_{i=1}^n \in \mathbb{R}^n$, then

$$\frac{d\rho_i}{dt} + \frac{1}{\Delta x^2} \sum_{j \in N(i)} g_{ij}(\rho) (F_i(\rho) - F_j(\rho)) = 0$$

holds for all $i \in V$. From the definition of $g_{ij}(\rho)$ in (4), we have (3). \square

To summarize, we have introduced a new discretization, which can be formally represented as

$$\frac{d\rho}{dt} = \operatorname{div}_G(\rho \nabla_G F(\rho)), \quad F(\rho) = \left(\frac{\partial}{\partial \rho_i} \mathcal{F}(\rho) \right)_{i=1}^n,$$

where

$$\operatorname{div}_G(\rho \nabla_G F(\rho)) = \frac{1}{\Delta x^2} \left(\sum_{j \in N(i)} \rho_j (F_j(\rho) - F_i(\rho))_+ - \sum_{j \in N(i)} \rho_i (F_i(\rho) - F_j(\rho))_+ \right)_{i=1}^n.$$

2.3. Gradient flow properties. Here, we show that (3) is a well defined ODE with gradient flow structures.

Theorem 5. *For any initial condition $\rho^0 \in \mathcal{P}_o(G)$, (3) has a unique solution $\rho(t) : [0, \infty) \rightarrow \mathcal{P}_o(G)$. Moreover,*

- (i) *there exists a constant $c = c(\rho^0) > 0$ depending on ρ^0 , such that $\rho_i(t) \geq c$ for all $i \in V$ and $t > 0$;*
- (ii) *the free energy $\mathcal{F}(\rho)$ is a Lyapunov function of (3):*

$$\frac{d}{dt} \mathcal{F}(\rho(t)) = - \sum_{(i,j) \in E} \left(\frac{F_i(\rho) - F_j(\rho)}{\Delta x} \right)_+^2 \rho_i.$$

Further, if $\lim_{t \rightarrow \infty} \rho(t)$ exists, call it ρ^∞ , then ρ^∞ is a Gibbs measure.

Proof. The proof of (i) can be found in [16], which is just a slight modification of [12]. Below, we only show (ii), which justifies saying that (3) is a gradient system. Firstly, we show that $\mathcal{F}(\rho)$ is a Lyapunov function:

$$\begin{aligned} \frac{d}{dt} \mathcal{F}(\rho(t)) &= \sum_{i=1}^n F_i(\rho) \cdot \frac{d\rho_i}{dt} \\ &= \frac{1}{\Delta x^2} \left\{ \sum_{i=1}^n \sum_{j \in N(i)} F_i(\rho)(F_j(\rho) - F_i(\rho))_+ \rho_j - \sum_{i=1}^n \sum_{j \in N(i)} F_i(\rho)(F_i(\rho) - F_j(\rho))_+ \rho_i \right\} \\ &\quad \text{Switch } i, j \text{ on the first formula} \\ &= \frac{1}{\Delta x^2} \left\{ \sum_{i=1}^n \sum_{j \in N(i)} F_j(\rho)(F_i(\rho) - F_j(\rho))_+ \rho_i - \sum_{i=1}^n \sum_{j \in N(i)} F_i(\rho)(F_i(\rho) - F_j(\rho))_+ \rho_i \right\} \\ &= - \sum_{(i,j) \in E} \left(\frac{F_i(\rho) - F_j(\rho)}{\Delta x} \right)_+^2 \rho_i \leq 0. \end{aligned}$$

Secondly, we prove that if $\rho^\infty = \lim_{t \rightarrow \infty} \rho(t)$ exists, then ρ^∞ is a Gibbs measure. Since $\mathcal{F}(\rho)$ is bounded in $\mathcal{P}(G)$, then $\lim_{t \rightarrow \infty} \frac{d}{dt} \mathcal{F}(\rho(t)) = 0$. From (i), we know that $\rho_i^\infty \geq c(\rho^0) > 0$ for any $i \in V$; so, the relation

$$\sum_{i=1}^n \sum_{j \in N(i)} (F_i(\rho^\infty) - F_j(\rho^\infty))_+^2 \rho_i^\infty = 0$$

implies $F_i(\rho^\infty) = F_j(\rho^\infty)$ for any $(i, j) \in E$. Since the graph is strongly connected,

$$F_i(\rho^\infty) = F_j(\rho^\infty), \quad \text{for any } i, j \in V.$$

Let

$$C := v_i + \sum_{j=1}^n w_{ij} \rho_j^\infty + \beta \log \rho_i^\infty, \quad \text{which is constant for any } i \in V,$$

$K = e^{-\frac{C}{\beta}}$ and use the fact $\sum_{i=1}^n \rho_i^\infty = 1$. Then, we have

$$\rho_i^\infty = \frac{1}{K} e^{-\frac{v_i + \sum_{j=1}^n w_{ij} \rho_j^\infty}{\beta}}, \quad K = \sum_{j=1}^n e^{-\frac{v_j + \sum_{i=1}^n w_{ij} \rho_i^\infty}{\beta}}.$$

Hence ρ^∞ is a Gibbs measure, which finishes the proof. \square

Notice that $(\mathcal{P}_o(G), W_2)$ is not a smooth Riemannian manifold, since for fixed i and $j \in V$, $g_{ij}(\rho)$ may be discontinuous with respect to ρ . Still, even though $(\mathcal{P}_o(G), W_2)$ is not smooth, (3) is a well defined ODE for any initial condition $\rho^0 \in \mathcal{P}_o(G)$.

One may be surprised by the unusual discretization of the Laplacian term, namely

$$\frac{1}{\Delta x^2} (\log \rho_j - \log \rho_i) g_{ij}(\rho) \tag{9}$$

which is different from the commonly adopted centered difference. We call (9) the ‘‘Log-Laplacian.’’ We observe that the Log-Laplacian plays a crucial role in the spatial discretization. Not only it implies that (3)’s equilibria are Gibbs measures, but it also indicates

that the boundary of the probability set⁸, $\partial\mathcal{P}(G)$, is a repeller for (3), which coincides with what is known in continuous states. We will see that this boundary repeller property plays a key role in the convergence result of section 3.

3. DISSIPATION RATE TO A DISCRETE GIBBS MEASURE

Considering the gradient flow (3), an important question arises. Assuming that $\rho(t)$ converges to an equilibrium ρ^∞ , how fast is the convergence speed? In the sequel, we show that the rate of convergence is exponential. Indeed, we capture such rate by the following explicit formula.

Definition 6. Denote

$$f_{ij} = \frac{\partial^2}{\partial \rho_i \partial \rho_j} \mathcal{F}(\rho),$$

and

$$h_{ij,kl} = \frac{f_{ik} + f_{jl} - f_{il} - f_{jk}}{\Delta x^2} \quad \text{for any } i, j, k, l \in V.$$

We define

$$\lambda_{\mathcal{F}}(\rho) = \min_{(\Phi_i)_{i=1}^n \in D} \sum_{(i,j) \in E} \sum_{(k,l) \in E} h_{ij,kl} \left(\frac{\Phi_i - \Phi_j}{\Delta x} \right)_+ \rho_i \left(\frac{\Phi_k - \Phi_l}{\Delta x} \right)_+ \rho_k,$$

where

$$D = \{(\Phi_i)_{i=1}^n \in \mathbb{R}^n \mid \sum_{(i,j) \in E} \left(\frac{\Phi_i - \Phi_j}{\Delta x} \right)_+^2 \rho_i = 1\}.$$

Remark 1. $\lambda_{\mathcal{F}}$ in Definition 6 plays the role of the smallest eigenvalue of the Hessian operator (on metric manifold) of the free energy at Gibbs measure; see [16] for more details about this connection.

Based on $\lambda_{\mathcal{F}}(\rho)$, we show the exponential convergence result for (3). We will assume that ρ^0 is in the basin of attraction of ρ^∞ for the gradient flow. I.e., if $\rho(t)$ is a solution of (3) with initial condition ρ^0 , then

$$(A) \quad \lim_{t \rightarrow \infty} \rho(t) = \rho^\infty \quad \text{and} \quad \rho^\infty \text{ is an isolated equilibrium.}$$

Theorem 7. Let (A) hold, and let $\lambda_{\mathcal{F}}(\rho^\infty) > 0$. Then there exists a constant $C = C(\rho^0, G) > 0$, depending on ρ^0 and G , such that

$$\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty) \leq e^{-Ct} (\mathcal{F}(\rho^0) - \mathcal{F}(\rho^\infty)).$$

Moreover, the asymptotic convergence rate is $2\lambda_{\mathcal{F}}(\rho^\infty)$. I.e., for any sufficiently small $\epsilon > 0$, there exists a time $T > 0$ depending on ϵ and ρ , such that when $t > T$,

$$\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty) \leq e^{-2(\lambda_{\mathcal{F}}(\rho^\infty) - \epsilon)t} (\mathcal{F}(\rho(T)) - \mathcal{F}(\rho^\infty)).$$

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$\partial\mathcal{P}(G) = \{(\rho_i)_{i=1}^n \mid \sum_{i=1}^n \rho_i = 1, \text{ there exists some index } i, \text{ such that } \rho_i = 0\}.$

Motivation of the proof. Our proof is motivated by some known facts of gradient flows in \mathbb{R}^n . We consider a λ -convex energy $g(x) \in C^2(\mathbb{R}^n)$ ⁹. The gradient flow associated to g is

$$\frac{dx_t}{dt} = -\nabla g(x_t), \quad x_t \in \mathbb{R}^n.$$

We compare the first and second derivative of $g(x_t)$ with respect to t :

$$\begin{aligned} \frac{d}{dt}g(x_t) &= -(\nabla g(x_t), \nabla g(x_t)), \\ \frac{d^2}{dt^2}g(x_t) &= -2(\text{Hess}_{\mathbb{R}^n}g(x_t) \cdot \nabla g(x_t), \nabla g(x_t)) \geq -2\lambda \frac{d}{dt}g(x_t). \end{aligned}$$

From the above comparison, we obtain the convergence result. Integrating on the time interval $[t, +\infty)$,

$$\frac{d}{dt}[g(x_t) - g(x_\infty)] \leq -2\lambda[g(x_t) - g(x_\infty)],$$

and applying Gronwall's inequality, the energy function $g(x_t)$ decreases exponentially

$$g(x_t) - g(x_\infty) \leq e^{-2\lambda t}(g(x_0) - g(x_\infty)).$$

In addition, from the dynamical viewpoint, the strict convexity of the free energy can be weakened: if the equilibrium x_∞ is a strict local minimizer, the exponential convergence result is still valid. Furthermore, the asymptotic convergence rate is $2\lambda_{\min} \text{Hess}_{\mathbb{R}^n}g(x_\infty)$ ¹⁰. \square

Proof of Theorem 7. Motivated by the standard approach in \mathbb{R}^n , we briefly sketch our proof (see [16] for complete details). The main idea is to compare the first and second derivatives of the free energy along (3).

Claim:

$$\begin{aligned} \frac{d^2}{dt^2}\mathcal{F}(\rho(t)) &= \frac{1}{\Delta x^4} \sum_{(i,j) \in E} \sum_{(k,l) \in E} h_{ij,kl}(F_i - F_j)_+ \rho_i(F_k - F_l)_+ \rho_k \\ &\quad + o\left(\frac{d}{dt}\mathcal{F}(\rho(t))\right). \end{aligned} \tag{10}$$

Here we denote $\lim_{h \rightarrow 0} \frac{o(h)}{h} = 0$, $F_i = \frac{\partial}{\partial \rho_i} \mathcal{F}(\rho)$, $f_{ij} = \frac{\partial^2}{\partial \rho_i \partial \rho_j} \mathcal{F}(\rho)$ and $h_{ij,kl} = f_{ik} + f_{lj} - f_{il} - f_{jk}$. If (10) holds, it is not hard to check that Theorem 7 holds.

Let's show (10) directly. Recall the gradient flow (3)

$$\frac{d\rho_i}{dt} = \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} (F_j - F_i)_+ \rho_j - \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \right\}.$$

We compute the first derivative of the free energy along (3):

$$\frac{d}{dt}\mathcal{F}(\rho(t)) = - \sum_{i=1}^n \sum_{j \in N(i)} \left(\frac{F_i - F_j}{\Delta x} \right)_+^2 \rho_i.$$

⁹ $\text{Hess}_{\mathbb{R}^n}g(x) \succeq \lambda I$, $\lambda > 0$ for all $x \in \mathbb{R}^n$.

¹⁰The smallest eigenvalue of the Hessian matrix of the energy at the minimizer.

Then we obtain the second derivative¹¹ by using the product rule:

$$\begin{aligned} \frac{d^2}{dt^2} \mathcal{F}(\rho(t)) = & - \sum_{i=1}^n \sum_{j \in N(i)} \left(\frac{F_i - F_j}{\Delta x} \right)_+^2 \frac{d\rho_i}{dt} \\ & - 2 \frac{1}{\Delta x^2} \sum_{i=1}^n \sum_{j \in N(i)} \left(\frac{dF_i}{dt} - \frac{dF_j}{dt} \right) (F_i - F_j)_+ \rho_i. \end{aligned} \quad \clubsuit \quad \spadesuit$$

Hence, (10) can be shown by the following two steps. Firstly, since $\rho(t)$ is assumed to converge to an equilibrium ρ^∞ and the boundary is a repeller (Theorem 5), we know that $\frac{d\rho}{dt} \rightarrow 0$ while $\rho_i(t) \geq c(\rho^0) > 0$. Hence \clubsuit is a high order term of the first derivative

$$\clubsuit = o\left(\frac{d}{dt} \mathcal{F}(\rho(t))\right).$$

Secondly, we have the following Lemma.

Lemma 8.

$$\spadesuit = 2 \sum_{(i,j) \in E} \sum_{(k,l) \in E} h_{ij,kl} \left(\frac{F_i - F_j}{\Delta x} \right)_+ \rho_i \left(\frac{F_k - F_l}{\Delta x} \right)_+ \rho_k.$$

¹¹Notice that $\frac{d^2}{dt^2} \mathcal{F}(\rho(t))$ exists for all $t \geq 0$, because $(F_i(\rho) - F_j(\rho))_+^2$ is differentiable everywhere with respect to ρ .

Proof of Lemma 8. We derive this result by a direct calculation¹².

$$\begin{aligned}
-\frac{1}{2} \spadesuit &= \frac{1}{\Delta x^2} \sum_{i=1}^n \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \left(\frac{d}{dt} F_i(\rho(t)) - \frac{d}{dt} F_j(\rho(t)) \right) \\
&= \frac{1}{\Delta x^2} \sum_{i=1}^n \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \left(\sum_{k=1}^n \frac{\partial F_i}{\partial \rho_k} \frac{d \rho_k}{dt} - \sum_{k=1}^n \frac{\partial F_j}{\partial \rho_k} \frac{d \rho_k}{dt} \right) \\
&= \frac{1}{\Delta x^2} \sum_{i=1}^n \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \sum_{k=1}^n (f_{ik} - f_{kj}) \frac{d \rho_k}{dt} \\
&= \frac{1}{\Delta x^4} \sum_{i=1}^n \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \sum_{k=1}^n (f_{ik} - f_{kj}) \left[\sum_{l \in N(k)} (F_l - F_k)_+ \rho_l - \sum_{l \in N(k)} (F_k - F_l)_+ \rho_l \right] \\
&= \frac{1}{\Delta x^4} \sum_{i=1}^n \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \left\{ \sum_{k=1}^n \sum_{l \in N(k)} (f_{ik} - f_{kj})(F_l - F_k)_+ \rho_l \right. \\
&\quad \left. - \sum_{k=1}^n \sum_{l \in N(k)} (f_{ik} - f_{kj})(F_k - F_l)_+ \rho_k \right\}
\end{aligned}$$

Relabel k, l in the first formula

$$\begin{aligned}
&= \frac{1}{\Delta x^4} \sum_{i=1}^n \sum_{j \in N(i)} (F_i - F_j)_+ \rho_i \left\{ \sum_{k=1}^n \sum_{l \in N(k)} (f_{il} - f_{lj})(F_k - F_l)_+ \rho_k \right. \\
&\quad \left. - \sum_{k=1}^n \sum_{l \in N(k)} (f_{ik} - f_{kj})(F_k - F_l)_+ \rho_k \right\} \\
&= \frac{1}{\Delta x^4} \sum_{i=1}^n \sum_{j \in N(i)} \sum_{k=1}^n \sum_{l \in N(k)} (f_{il} - f_{lj} - f_{ik} + f_{kj})(F_i - F_j)_+ \rho_i (F_k - F_l)_+ \rho_k \\
&= \frac{1}{\Delta x^4} \sum_{(i,j) \in E} \sum_{(k,l) \in E} (f_{il} - f_{lj} - f_{ik} + f_{kj})(F_i - F_j)_+ \rho_i (F_k - F_l)_+ \rho_k.
\end{aligned}$$

□

Combining all the above facts, the claim and the proof of Theorem 7 follow. □

3.1. Analysis of dissipation rate. In the sequel, we further elucidate the relationship between convexity of the free energy (Hessian operator in \mathbb{R}^n) and the dissipation rate.

Lemma 9. *Denote*

$$\tilde{div}_G(\rho \nabla_G \Phi) := \left(\frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} (\Phi_i - \Phi_j)_+ \rho_i - \sum_{j \in N(i)} (\Phi_j - \Phi_i)_+ \rho_j \right\} \right)_{i=1}^n.$$

¹²Here we use the relabeling technique heavily: For a matrix $(k_{ij})_{1 \leq i, j \leq n}$,

$$\sum_{(i,j) \in E} k_{ij} = \sum_{(j,i) \in E} k_{ji}.$$

Then $\lambda_{\mathcal{F}}(\rho)$ in Definition 6 is equivalent to

$$\lambda_{\mathcal{F}}(\rho) = \min\left\{\left(\tilde{\operatorname{div}}_G(\rho \nabla_G \Phi)\right)^T \operatorname{Hess}_{\mathbb{R}^n} \mathcal{F}(\rho) \tilde{\operatorname{div}}_G(\rho \nabla_G \Phi) : \sum_{(i,j) \in E} \left(\frac{\Phi_i - \Phi_j}{\Delta x}\right)_+^2 \rho_i = 1\right\}.$$

The proof of Lemma 9 is based on a direct computation (see details in [16]).

Lemma 9 gives convergence rates for many semi-discretization schemes.

Corollary 10. *Consider the gradient flow (3) of the free energy*

$$\mathcal{F}(\rho) = \sum_{i=1}^n v_i \rho_i + \frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n w_{ij} \rho_i \rho_j + \beta \sum_{i=1}^n \rho_i \log \rho_i.$$

If the matrix $W = (w_{ij})_{1 \leq i,j \leq n}$ is semi positive definite, then there is a unique Gibbs measure ρ^∞ , which is a global attractor of (3). Moreover, there exists a constant $C > 0$, such that

$$\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty) \leq e^{-Ct} (\mathcal{F}(\rho^0) - \mathcal{F}(\rho^\infty))$$

with asymptotic rate $2\lambda_{\mathcal{F}}(\rho^\infty)$.

Proof. The main idea of proof is as follows (full details are in [?]). Notice that since $\operatorname{Hess}_{\mathbb{R}^n} \sum_{i=1}^n \rho_i \log \rho_i = \operatorname{diag}(\frac{1}{\rho_k})_{1 \leq k \leq n}$ and the matrix W is semi positive definite, then

$$\operatorname{Hess}_{\mathbb{R}^n} \mathcal{F}(\rho)|_{\rho=\rho^\infty} = W + \beta \operatorname{diag}(\frac{1}{\rho_k^\infty})_{1 \leq k \leq n}$$

is a positive definite matrix. Then, from Lemma 9 and Theorem 7, we know that (3) converges exponentially. \square

Throughout this section, we observe another important effect of the Log-Laplacian, which reflects the convexity property of the linear entropy

$$\mathcal{H}(\rho) = \sum_{i=1}^n \rho_i \log \rho_i.$$

Lemma 9 says that

$$\lambda_{\mathcal{H}}(\rho) = \min\left\{\sum_{i=1}^n \frac{1}{\rho_i} (\tilde{\operatorname{div}}_G(\rho \nabla_G \Phi)|_i)^2 : \sum_{(i,j) \in E} \left(\frac{\Phi_i - \Phi_j}{\Delta x}\right)_+^2 \rho_i = 1\right\}.$$

Given any Gibbs measure ρ^∞ , we know that $\lambda_{\mathcal{H}}(\rho^\infty) > 0$. To visualize that, consider a simple example with no interaction energy, meaning that $(w_{ij}) = 0$. In this case, (3) is a semi-discretization for a linear Fokker-Planck equation. The free energy is

$$\mathcal{F}(\rho) = \sum_{i=1}^n v_i \rho_i + \beta \mathcal{H}(\rho).$$

Here, strict convexity of $\mathcal{H}(\rho)$ tells that there always exists a constant $C > 0$, such that

$$\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty) \leq e^{-Ct} (\mathcal{F}(\rho^0) - \mathcal{F}(\rho^\infty))$$

holds with asymptotic rate $2\lambda_{\mathcal{H}}(\rho^\infty)$.

4. NUMERICAL ANALYSIS

In this section, we show some numerical properties of (3).

4.1. Spatial consistency. To begin with, we show that (3) is a finite volume scheme for the PDE (1). For concreteness, we use a lattice graph. Rewrite (3) in the following form

$$\frac{d\rho_i}{dt} = \frac{1}{\Delta x^2} \left\{ \sum_{v=1}^d \sum_{j \in N_v(i)} [F_j(\rho) - F_i(\rho)]_+ \rho_j - \sum_{v=1}^d \sum_{j \in N_v(i)} [F_i(\rho) - F_j(\rho)]_+ \rho_i \right\}.$$

Denote $i = (i_1, \dots, i_d)$, and G is a cartesian graph of d one dimensional lattices, i.e. $G = G_1 \square \dots \square G_d$ with $G_v = (V_v, E_v)$. Here

$$N_v(i) = \{(i_1, \dots, i_{v-1}, j_v, i_{v+1}, \dots, i_d) \in V \mid (i_v, j_v) \in E_v\}.$$

Theorem 11. *The semi-discretization (3) is a consistent finite volume scheme for the PDE (1).*

Proof. Denote by $\rho_i(t)$ a discrete probability function

$$\rho_i(t) = \int_{C_i} \rho(t, x) dx,$$

where C_i is a cube in \mathbb{R}^d centered at point i with equal width Δx . Here $i \in V$ represents a point $x(i) \in \mathbb{R}^d$. Let $e_v = (0, \dots, 1, \dots, 0)^T$, where 1 is in the v -th position, $v = 1, \dots, d$. So in this setting, $N_v(i)$ for a lattice graph only contains the two points $x(i) - e_v \Delta x$, $x(i) + e_v \Delta x$. Denote $\rho_j(t)$ by

$$\rho_j(t) = \int_{C_{i_+}} \rho(t, x) dx,$$

where $j \in N(i)$ satisfies $x(j) = x(i) + e_v \Delta x$ and C_{i_+} is a cube centered at the point $j \in V$.

Without loss of generality, we assume $F(x(i) + e_v \Delta x, \rho) \geq F(x(i), \rho) \geq F(x(i) - e_v \Delta x, \rho)$. Applying Taylor expansion of (3) relative to the direction e_v , we obtain

$$\begin{aligned}
& \frac{1}{\Delta x^2} \left\{ \sum_{j \in N_v(i)} [F_j(\rho) - F_i(\rho)]_{+\rho_j} - \sum_{j \in N_v(i)} [F_i(\rho) - F_j(\rho)]_{+\rho_j} \right\} \\
&= \frac{1}{\Delta x^2} \left\{ [F(x(i) + e_v \Delta x, \rho) - F(x(i), \rho)] \int_{C_{i+}} \rho(t, x) dx \right. \\
&\quad \left. - [F(x(i), \rho) - F(x(i) - e_v \Delta x, \rho)] \int_{C_i} \rho(t, x) dx \right\} \\
&= \frac{1}{\Delta x^2} \left\{ \left[\frac{\partial F}{\partial x_v}(x(i), \rho) \Delta x + \frac{1}{2} \frac{\partial F}{\partial x_v}(x(i), \rho) \Delta x^2 \right] \int_{C_{i+}} \rho(t, x) dx \right. \\
&\quad \left. - \left[\frac{\partial F}{\partial x_v}(x(i), \rho) \Delta x - \frac{1}{2} \frac{\partial}{\partial x_v} F(x(i), \rho) \Delta x^2 \right] \int_{C_i} \rho(t, x) dx + O(\Delta x^3) \right\} \\
&= \frac{1}{\Delta x} \frac{\partial F}{\partial x_v}(x(i), \rho) \left[\int_{C_{i+}} \rho(t, x) dx - \int_{C_i} \rho(t, x) dx \right] \\
&\quad + \frac{1}{2} \frac{\partial^2 F}{\partial x_v^2}(x(i), \rho) \left[\int_{C_{i+}} \rho(t, x) dx + \int_{C_i} \rho(t, x) dx \right] + O(\Delta x) \\
&= \frac{\partial F}{\partial x_v}(x(i), \rho) \int_{C_i} \frac{\rho(t, x + e_v \Delta x) - \rho(t, x)}{\Delta x} dx \\
&\quad + \frac{\partial^2 F}{\partial x_v^2}(x(i), \rho) \int_{C_i} \frac{\rho(t, x + e_v \Delta x) + \rho(t, x)}{2} dx + O(\Delta x) \\
&= \int_{C_i} \nabla_{x_v} \cdot (\rho(t, x) \nabla_{x_v} F(x, \rho)) dx + O(\Delta x).
\end{aligned} \tag{11}$$

Similarly, we can show the same results for other possible configurations, such as $F(x(i) - e_v \Delta x, \rho) \geq F(x(i), \rho) \geq F(x(i) + e_v \Delta x, \rho)$, $F(x(i), \rho) \geq F(x_v - e_v \Delta x, \rho) \geq F(x(i) + e_v \Delta x, \rho)$.

Therefore, combining all directions e_v with $v = 1, \dots, d$, the right-hand-side of (3) becomes

$$\begin{aligned}
& \frac{d\rho_i}{dt} - \frac{1}{\Delta x^2} \sum_{v=1}^d \left\{ \sum_{j \in N_v(i)} [F_j(\rho) - F_i(\rho)]_{+\rho_j} - \sum_{j \in N_v(i)} [F_i(\rho) - F_j(\rho)]_{+\rho_j} \right\} \\
&= \int_{C_i} \left\{ \frac{\partial \rho(t, x)}{\partial t} - \sum_{v=1}^d \nabla_{x_v} \cdot (\rho(t, x) \nabla_{x_v} F(x, \rho)) \right\} dx + dO(\Delta x) \\
&= \int_{C_i} \left\{ \frac{\partial \rho(t, x)}{\partial t} - \nabla \cdot (\rho(t, x) \nabla_x F(x, \rho)) \right\} dx + dO(\Delta x) \\
&= O(\Delta x).
\end{aligned}$$

This shows that (3) is a finite volume first order semi-discretization scheme for (1). \square

Remark 2. Although we proved consistency of the scheme by using a square lattice graph, we can easily extend it to other regular polygonal discretizations, e.g. equilateral triangular grids in \mathbb{R}^2 .

4.2. Time discretization. To deal with the time discretization, we use a forward Euler scheme on (3):

$$\frac{\rho_i^{k+1} - \rho_i^k}{\Delta t} = \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j^k (F_j(\rho^k) - F_i(\rho^k))_+ - \sum_{j \in N(i)} \rho_i^k (F_i(\rho^k) - F_j(\rho^k))_+ \right\}. \quad (12)$$

Lemma 12. *Assume that the discrete free energy $\mathcal{F}(\rho)$ is strictly convex on $\mathcal{P}_o(G)$.*

(i) *For a given small tolerance constant $\epsilon > 0$, and initial measure $\rho^0 \in \mathcal{P}_o(G)$, there exists a finite time $T = O(\log \frac{1}{\epsilon})$, such that when $t > T$,*

$$|\mathcal{F}(\rho(t)) - \mathcal{F}(\rho^\infty)| < \epsilon.$$

(ii) *There exists a constant h , such that if $0 < \Delta t \leq h$, $\rho^k = (\rho_i^k)_{i=1}^n \in \mathcal{P}_o(G)$, for all $k = 0, 1, \dots, [\frac{T}{\Delta t}]$, where T is the value from (i).*

Proof. (i) can be shown by the exponential convergence result in Corollary 10. Since there exists a constant $C > 0$, such that

$$\mathcal{F}(\rho(T)) - \mathcal{F}(\rho^\infty) \leq e^{-CT} (\mathcal{F}(\rho^0) - \mathcal{F}(\rho^\infty)),$$

then if $\rho(T)$ satisfies $|\mathcal{F}(\rho(T)) - \mathcal{F}(\rho^\infty)| < \epsilon$, we need to set

$$T \geq \frac{1}{C} \log \frac{\mathcal{F}(\rho^0) - \mathcal{F}(\rho^\infty)}{\mathcal{F}(\rho(T)) - \mathcal{F}(\rho^\infty)}.$$

In other words, we can approximate ρ^∞ with $O(\epsilon)$ precision by time $T = O(\log \frac{1}{\epsilon})$.

We prove (ii) in two steps.

Firstly, we show that $\rho^k = (\rho_i^k)_{i=1}^n$ stays positive ($\min_{i \in V} \rho_i^k > 0$) for all $k = 1, \dots, N$. From Theorem 4, we know that the boundary is a repeller for (3). This means that there exists a constant $\epsilon_0 = \epsilon_0(\rho^0) > 0$, such that

$$\min_{i \in \{1, \dots, n\}} \rho_i(t) \geq \epsilon_0(\rho^0), \quad \text{for all } t \geq 0.$$

Since the forward Euler scheme is convergent for Lipschitz right-hand-sides (and this is the case for us), there exists constant h , such that when $\Delta t \leq h$, we have

$$\min_{i \in \{1, \dots, n\}} |\rho_i(k\Delta t) - \rho_i^k| \leq \frac{\epsilon_0}{2},$$

from which $\min_{i \in \{1, \dots, n\}} \rho_i^k \geq \frac{1}{2}\epsilon_0 > 0$.

Secondly, we show that $\sum_{i=1}^n \rho_i^k = 1$ for all $k = 1, \dots, N$. Since $\sum_{i=1}^n \rho^0 = 1$, it is sufficient to prove that

$$\sum_{i=1}^n \rho_i^{k+1} = \sum_{i=1}^n \rho_i^k, \quad \text{for any } k.$$

This is a linear invariant, and it is therefore kept by Euler method. Indeed, an explicit computation gives

$$\begin{aligned}
& \sum_{i=1}^n \frac{\rho_i^{k+1} - \rho_i^k}{\Delta t} \\
&= \sum_{i=1}^n \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j^k (F_j(\rho^k) - F_i(\rho^k))_+ - \sum_{j \in N(i)} \rho_i^k (F_i(\rho^k) - F_j(\rho^k))_+ \right\} \\
&= \sum_{i=1}^n \sum_{j \in N(i)} (F_j(\rho^k) - F_i(\rho^k)) g_{ij}(\rho^k) \\
&= \sum_{(i,j) \in E} F_j(\rho^k) g_{ij}(\rho^k) - \sum_{(i,j) \in E} F_i(\rho^k) g_{ij}(\rho^k) \\
&\quad \text{Relabel } i \text{ and } j \text{ for the first formula and notice } g_{ij} = g_{ji} \\
&= \sum_{(i,j) \in E} F_i(\rho^k) g_{ji}(\rho^k) - \sum_{(i,j) \in E} F_i(\rho^k) g_{ij}(\rho^k) = 0.
\end{aligned}$$

□

Remark 3. In practice, cfr. with [8], we may consider¹³ $\Delta t \leq \frac{\Delta x^2}{\Delta(G)M}$, with $M = 2 \sup_{i \in V} |F_i(\rho^k)|$ and $\Delta(G)$ representing the maximal degree¹⁴ of the graph G . For sufficiently small Δt , we know that M will be a bounded function up to a finite time T .

4.3. An extension. We extend the idea of semi-discretization scheme (3) to deal with more general Fokker-Planck equations. Consider

$$\frac{\partial \rho}{\partial t} = \nabla \cdot [\rho(f_v(x, \rho))_{v=1}^d]. \quad (13)$$

Here, (13) may fail to be a gradient flow with respect to the 2-Wasserstein metric. In this case, we cannot consider a discretization which is a gradient flow of a certain free energy. However, we can still construct a flow (semi-discretization scheme) whose solutions lie on the probability set. The observation to use is that there always exists functions $(u_v(x, \rho))_{v=1}^d$ such that

$$\nabla_{x_v} u_v(x, \rho) = f_v(x, \rho), \quad \text{for } v \in \{1, \dots, d\}.$$

Example 1 (van der Pol). Consider the 2 dimensional Fokker-Planck equation

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \begin{pmatrix} x_2 \\ \alpha(1 - x_1^2) - x_2 \end{pmatrix}) + \frac{\partial^2 \rho}{\partial x_2^2} = \nabla \cdot (\rho \begin{pmatrix} -f_1(x) \\ -f_2(x) \end{pmatrix}),$$

where $x = (x_1, x_2)$, $f_1(x, \rho) = x_2$ and $f_2(x, \rho) = (1 - x_1^2) - x_2 + \nabla_{x_2} \log \rho(x)$. We let

$$u_1(x) = \int f_1(x, \rho) dx_1 = x_1 x_2,$$

and

$$u_2(x) = \int f_2(x, \rho) dx_2 = (1 - x_1^2) x_2 - \frac{1}{2} x_2^2 + \log \rho(x_1, x_2).$$

¹³This is a classical CFL condition for a lattice graph.

¹⁴The number of edges incident to any vertex $i \in V$.

Then the Fokker-Planck equation becomes

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \begin{pmatrix} -\nabla_{x_1} u_1(x) \\ -\nabla_{x_2} u_2(x) \end{pmatrix}).$$

Based on the above observation, we naturally extend (3) to the semi-discretization of (13)

$$\frac{d\rho_i}{dt} = \frac{1}{\Delta x^2} \left\{ \sum_{v=1}^d \sum_{j \in N_v(i)} [u_v(j, \rho) - u_v(i, \rho)]_+ \rho_j - \sum_{v=1}^d \sum_{j \in N_v(i)} [u_v(i, \rho) - u_v(j, \rho)]_+ \rho_i \right\}. \quad (14)$$

We observe that (3) is a special case of (14). Similarly to Theorem 11, we can show that the semi-discretization (14) is a consistent finite volume scheme for (13).

5. NUMERICAL EXPERIMENTS

In this section, we illustrate the proposed semi-discretization with several numerical experiments.

Example 2 (Nonlinear Fokker-Planck equation). *We consider a nonlinear interaction-diffusion equation in granular gas [4, 23],*

$$\frac{\partial \rho}{\partial t} = \nabla \cdot [\rho \nabla (W * \rho + V(x))] + \beta \Delta \rho,$$

where $W(x, y) = \frac{1}{3} \|x - y\|^3$ and $V(x) = \frac{\|x\|^2}{2}$ with $\|\cdot\|$ the 2 norm in \mathbb{R}^d , $d = 1, 2$.

The PDE has a unique stationary measure (Gibbs measure),

$$\rho^*(x) = \frac{1}{K} e^{-\frac{\int_{\mathbb{R}^d} W(x, y) \rho^*(y) dy + V(x)}{\beta}}, \quad \text{where } K = \int_{\mathbb{R}^d} e^{-\frac{\int_{\mathbb{R}^2} W(x, y) \rho^*(y) dy + V(x)}{\beta}} dx.$$

We apply (3) to discretize this PDE with $\beta = 0.01$:

$$\begin{aligned} \frac{d\rho_i}{dt} = & \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j \left(\sum_{i=1}^n w_{ij} \rho_i - \sum_{j=1}^n w_{ij} \rho_j + v_j - v_i + \beta \log \rho_j - \beta \log \rho_i \right) + \right. \\ & \left. - \sum_{j \in N(i)} \rho_i \left(\sum_{j=1}^n w_{ij} \rho_j - \sum_{i=1}^n w_{ij} \rho_i + v_i - v_j + \beta \log \rho_i - \beta \log \rho_j \right)_+ \right\}, \end{aligned}$$

and further discretize in time with the forward Euler method (12) with time step $\Delta t = 10^{-4}$ and initial condition $\rho_i^0 = \frac{1}{L} e^{-\frac{\|x(i)\|^2}{200}}$, $L = \sum_{i=1}^n e^{-\frac{\|x(i)\|^2}{200}}$.

When $d = 1$, we take a one dimensional lattice graph of $[-5, 5]$ with $\Delta x = 0.5$; see Figure 1

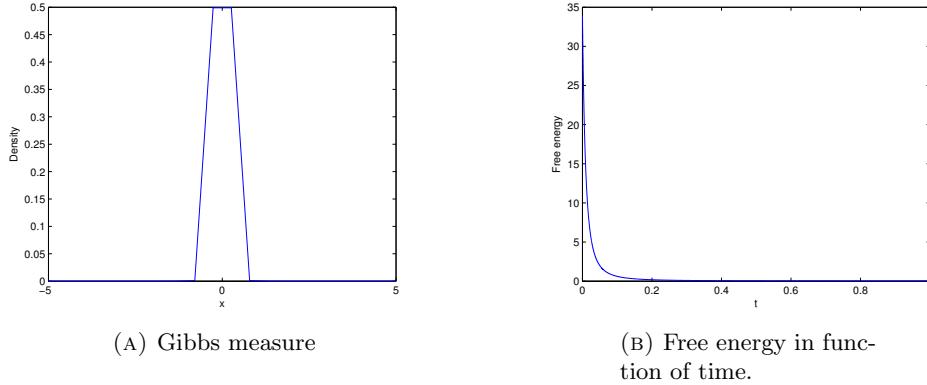


FIGURE 1. Example 2: 1-d.

When $d = 2$, we consider a two dimensional lattice graph of $[-5, 5] \times [-5, 5]$ with $\Delta x = 0.5$; see Figure 2

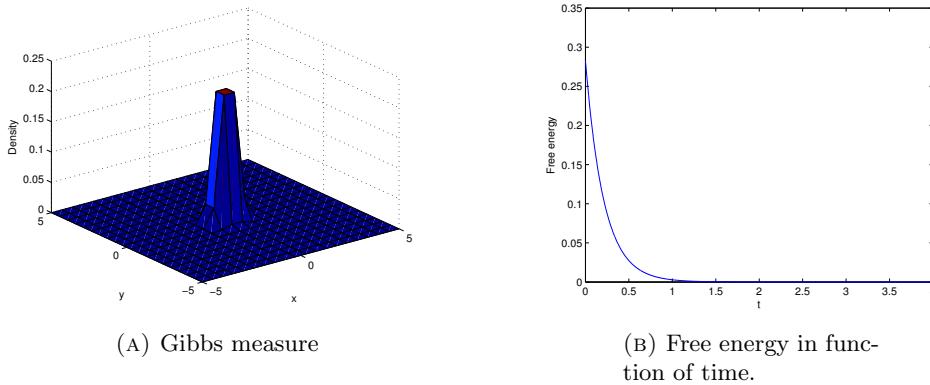


FIGURE 2. Example 2: 2-d.

It is known, see [10], that solutions of this PDE converge to the unique Gibbs measure, which itself converges to a δ -measure supported at the origin when $\beta \rightarrow 0$. In addition, the solution converges to the Gibbs measure exponentially. We observe that (3) reflects all of these behaviors and the free energy along solutions of (3) decreases exponentially.

Example 3 (Linear Fokker-Planck equation). We consider a linear Fokker-Planck equation

$$\frac{\partial \rho}{\partial t} = \nabla \cdot [\rho \nabla V(x)] + \beta \Delta \rho, \quad (15)$$

with a potential function $V(x) = \frac{\|x\|^4}{4} - \frac{\|x\|^2}{2}$. Here the underlying state is \mathbb{R}^d , $d = 1, 2$. In this case, the unique Gibbs measure is

$$\rho^*(x) = \frac{1}{K} e^{-\frac{V(x)}{\beta}}, \quad \text{where} \quad K = \int_{\mathbb{R}^d} e^{-\frac{V(x)}{\beta}} dx.$$

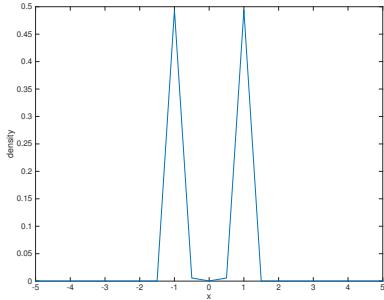
We use (3) to approximate the solution of this PDE with $\beta = 0.01$,

$$\begin{aligned} \frac{d\rho_i}{dt} = & \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j (v_j - v_i + \beta \log \rho_j - \beta \log \rho_i)_+ \right. \\ & \left. - \sum_{j \in N(i)} \rho_i (v_i - v_j + \beta \log \rho_i - \beta \log \rho_j)_+ \right\}, \end{aligned}$$

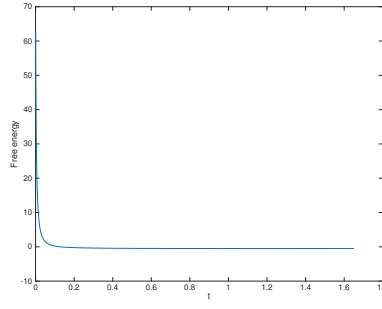
and further discretize in time by the forward Euler method (12) with time step $\Delta t = 10^{-4}$.

Initial condition is $\rho_i^0 = \frac{1}{L} e^{-\frac{\|x(i)\|^2}{200}}$, $L = \sum_{i=1}^n e^{-\frac{\|x(i)\|^2}{200}}$.

If $d = 1$, we take a uniform discretization of $[-5, 5]$ with $\Delta x = 0.5$; see Figure 3.



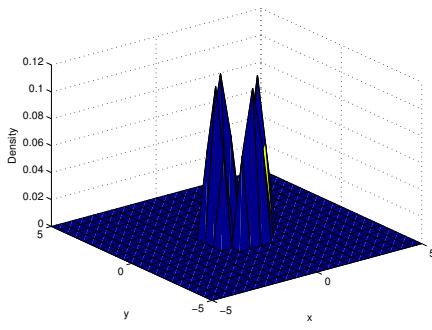
(A) Gibbs measure



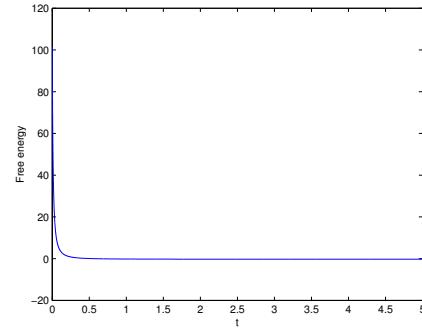
(B) Free energy with respect to time.

FIGURE 3. Example 3: 1-d.

If $d = 2$, we take a uniform discretization of $[-5, 5] \times [-5, 5]$ with $\Delta x = 0.5$; see Figure 4.



(A) Gibbs measure



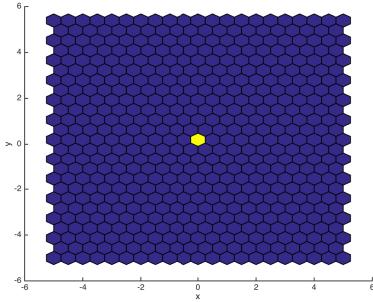
(B) Free energy with respect to time.

FIGURE 4. Example 3: 2-d.

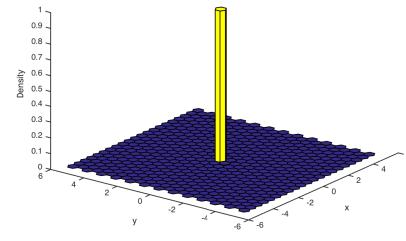
It is known, see [18], that the solution of the linear Fokker-Planck equation always converges to the Gibbs measure exponentially. The computational results in both cases reflects

this fact, in agreement with the discussion of Section 3. Note that here the potential function $V(x)$ is not strictly convex. It is the strict convexity of the entropy in probability set that plays the key role in convergence. This asymptotic convergence rate is fully determined by $\lambda_{\mathcal{H}}(\rho^\infty)$ in Definition 6.

Example 4 (Hexagonal graph). We apply (3) with a regular hexagonal graph to approximate the solution of the linear Fokker-Planck equation (15), where $\Delta x = \sqrt{3}a$ and $a = \frac{1}{2\sqrt{3}}$ is the length of the hexagon side. We consider the potential function $V(x) = \frac{\|x\|^2}{2}$ and $\beta = 0.01$ in \mathbb{R}^2 . We solve (3) similarly to Example 3.



(A) Hexagonal graph.



(B) Gibbs measure.

FIGURE 5. Example 4.

A node in the hexagonal graph typically has 6 neighbors, while in the square lattice just 4. The two graphs represent two different discretization of the spatial variable. However, in Theorem 5, we know that (3) always converges to its equilibrium, which is the discrete Gibbs measure of current spatial discretization.

Example 5 (General Fokker-Planck equation). We consider the Fokker-Planck equation [13]

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \left(\frac{x_2}{\alpha(1-x_1^2) - x_2} \right)) = \beta \Delta_{x_2} \rho,$$

whose underlying state is the stochastic van der Pol oscillator

$$\begin{aligned} dx_1 &= x_2 dt \\ dx_2 &= [\alpha(1-x_1^2)x_2 - x_1]dt + \sqrt{2\beta}dW_t. \end{aligned}$$

We apply the semi-discretization (14) to approximate the solution of this PDE,

$$\begin{aligned} \frac{d\rho_i}{dt} &= \frac{1}{\Delta x^2} \left\{ \sum_{j \in N_1(i)} \rho_j [u_{1j} - u_{1i}]_+ - \sum_{j \in N_1(i)} \rho_i [u_{1i} - u_{1j}]_+ \right. \\ &\quad \left. + \sum_{j \in N_2(i)} \rho_j [u_{2j} - u_{2i}]_+ - \sum_{j \in N_2(i)} \rho_i [u_{2i} - u_{2j}]_+ \right\}, \end{aligned}$$

where $i \in V \subset \mathbb{R}^2$ and

$$u_{1i} = -x_1 x_2|_{(x_1, x_2)=i}, \quad u_{2i} = -\alpha(1-x_1^2)x_2 + \frac{1}{2}x_2^2 + \beta \log \rho(x_1, x_2)|_{(x_1, x_2)=i}.$$

Further, we discretize in time by the forward Euler method (12) with time step $\Delta t = 10^{-4}$. Initial condition is $\rho_i^0 = \frac{1}{L} e^{-\frac{\|x(i)\|^2}{200}}$, $L = \sum_{i=1}^n e^{-\frac{\|x(i)\|^2}{200}}$.

Let $\alpha = 1$, $\beta = 0.125$, and consider a lattice graph on $[-10, 10] \times [-10, 10]$ with $\Delta x = 0.4$. The result in Figure 6 shows the obtained approximation of the stationary measure of the stochastic van der Pol oscillator.

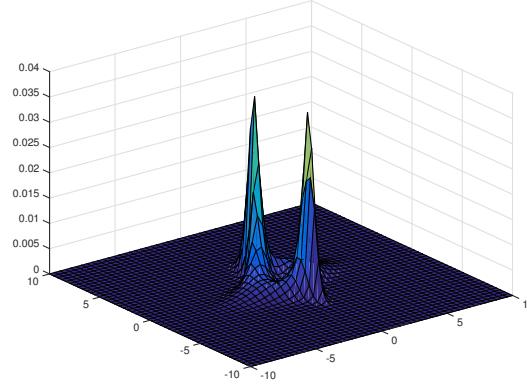


FIGURE 6. Example 5. Stationary measure, van der Pol.

Similarly, we consider the Fokker-Planck equation

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \begin{pmatrix} x_2 \\ -2\xi\omega x_2 + \omega x_1 - \omega^2 r x_1^3 \end{pmatrix}) = \beta \Delta_{x_2} \rho,$$

associated with the stochastic Duffing oscillator

$$\begin{aligned} dx_1 &= x_2 dt \\ dx_2 &= [-2\xi\omega x_2 + \omega x_1 - \omega^2 r x_1^3] dt + \sqrt{2\beta} dW_t. \end{aligned}$$

Let $\xi = 0.2$, $\omega = 1$, $r = 0.1$, $\beta = 0.125$ and a lattice graph of $[-10, 10] \times [-10, 10]$ with $\Delta x = 0.4$. The computed invariant measure is shown in Figure 7.

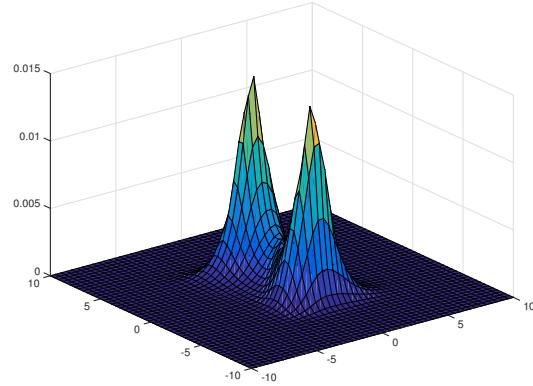


FIGURE 7. Example 5. Stationary measure, Duffing.

In these examples, we have shown that our discretization scheme (14) finds a two-peaks stationary measure, even though the underlying Fokker-Planck equations are not gradient flow type.

It is interesting to observe that, in the above two figures, stationary measures are supported around the limit cycles of the oscillators. The two peaks in the stationary measures reflect that there is slow and fast motion in the underlying dynamical systems; namely, the two peaks are witness to the fact that there is a larger probability that a trajectory at time t will be found in the slow motion region, see figure 8.

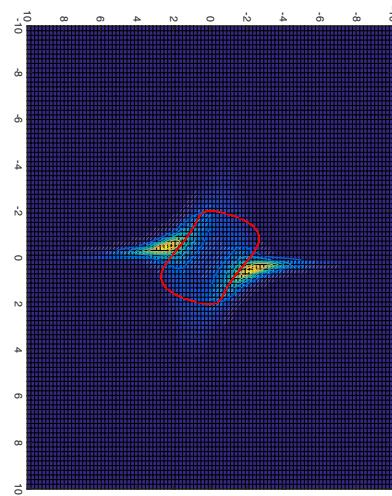


FIGURE 8. The plot of stationary measure and limit cycle (red) of van der Pol oscillator.

6. CONCLUSION

We have derived a new semi-discretization scheme (3) for the PDE (1). In comparison to other methods, our scheme (3) has the following advantages.

- (1) Firstly, our scheme (3) works on a finite graph, which is a spatial discretization of the underlying state. As a result of having this graph, we can handle a variety of boundary conditions, e.g. zero-flux conditions or periodic conditions, and different types of underlying states, such as \mathbb{R}^d , open set of \mathbb{R}^d , or Riemannian manifold.
- (2) Secondly, we derive (3) from the viewpoint of free energy and optimal transport. Hence, (3) can keep the gradient flow structure of (1). On one hand, this fact gives that (3) is a well defined flow whose equilibria are discrete Gibbs measures; on the other hand, solutions of (3) converge to a Gibbs measure with exponential rate. This property allowed us to discretize (3) in time by a forward Euler scheme.
- (3) Lastly, we bring a new twist to discretize the diffusion term, namely

$$\frac{1}{\Delta x^2} \sum_{j \in N(i)} (\log \rho_j - \log \rho_i) g_{ij}(\rho).$$

We called it Log-Laplacian, and it is quite different from commonly known centered differences or the Graph Laplacian. Although the log term brings some nonlinearities into the algorithm, it also brings many benefits. One is that solutions of (3) always stay in $\mathcal{P}_o(G)$, and thus remain positive and conserve the total probability automatically. The other is that the scheme naturally inherits the convexity of the entropy, a fact which plays a critical role in the convergence result.

Our results also open the door to many new research questions. “What is the asymptotic convergence rate of (3)?” This rate certainly depends on the boundary conditions of the PDE and on the explicit form of the interaction potential, and it deserves more investigation. In future works, we will study this problem and work on other gradient flow types of advection diffusion equations.

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7. APPENDIX

In general, it appears to be very difficult to obtain sharp information on the asymptotic convergence rates of Theorem 7, namely the quantity $\lambda_{\mathcal{F}}(\rho)$ there. Here, we consider a simple 1-d model problem, for which exact formulas for $\lambda_{\mathcal{F}}(\rho)$ can be given, and we can highlight the dependence of these convergence rates on the graph structure (hence the boundary conditions of the PDE).

A 1-d model problem. Suppose that the free energy contains only the linear entropy term, so that the gradient flow is the heat equation:

$$\frac{\partial \rho}{\partial t} = \Delta \rho, \quad x \in (a, b). \quad (16)$$

Here, we consider either (i) Neumann boundary conditions (zero flux) $\frac{\partial \rho}{\partial x}|_{x=a} = \frac{\partial \rho}{\partial x}|_{x=b} = 0$, or (ii) periodic boundary conditions $\rho(t, a) = \rho(t, b)$.

We approximate the solution of (16) by (3), with a uniform discretization $\Delta x = \frac{b-a}{n-1}$:

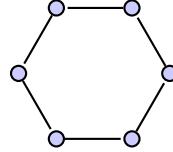
$$\frac{d\rho_i}{dt} = \frac{1}{\Delta x^2} \left\{ \sum_{j \in N(i)} \rho_j (\log \rho_j - \log \rho_i)_+ - \sum_{j \in N(i)} \rho_i (\log \rho_i - \log \rho_j)_+ \right\}. \quad (17)$$

The above two types of boundary conditions lead to distinct graph structures.

(i) A lattice graph L_n :



(ii) A cycle graph C_n :



In both cases, (17) is the gradient flow of the discrete linear entropy

$$\mathcal{H}(\rho) = \sum_{i=1}^n \rho_i \log \rho_i,$$

and the unique Gibbs measure is $\rho^\infty = (\frac{1}{n}, \dots, \frac{1}{n})$. We are going to estimate how fast the solution $\rho(t)$ of the semi-discretization scheme (17) converges to the equilibrium ρ^∞ .

As we have seen in Theorem 7, the asymptotic convergence rates are determined by $\lambda_{\mathcal{F}}(\rho)$:

$$\begin{aligned} \lambda_{\mathcal{H}}(\rho^\infty) = \min_{\Phi \in \mathbb{R}^n} \{ & \frac{1}{\Delta x^4} \sum_{(i,j) \in E} \sum_{(k,l) \in E} h_{ij,kl} (\Phi_i - \Phi_j)_+ (\Phi_k - \Phi_l)_+ : \\ & \sum_{(i,j) \in E} \left(\frac{\Phi_i - \Phi_j}{\Delta x} \right)_+^2 \rho_i = 1 \}, \end{aligned} \quad (18)$$

where

$$h_{ij,kl} = f_{ik} + f_{jl} - f_{il} - f_{jk}, \quad \text{and} \quad f_{ij}(\rho^\infty) = \frac{\partial^2}{\partial \rho_i \partial \rho_j} \mathcal{H}(\rho)|_{\rho=\rho^\infty} = \begin{cases} \frac{1}{\rho_i^\infty} & \text{if } i = j; \\ 0 & \text{if } i \neq j. \end{cases}$$

For the present model, we can find exact values of (18) for the above two graphs.

Theorem 13. *We have*

$$\lambda_{\mathcal{H}}(\rho^\infty) = \frac{\pi^2}{(b-a)^2} + o(1), \quad (L_n)$$

and

$$\lambda_{\mathcal{H}}(\rho^\infty) = \frac{4\pi^2}{(b-a)^2} + o(1). \quad (C_n)$$

Proof. First, consider the lattice graph L_n . Without loss of generality, let $(\Phi_i)_{i=1}^n$ in (18) satisfy the relation

$$\Phi_1 \geq \Phi_2 \geq \dots \geq \Phi_n. \quad (19)$$

Denote $\xi := (\xi_i)_{i=1}^{n-1} \in \mathbb{R}_+^{n-1}$ by

$$\xi_i := \frac{\Phi_{i+1} - \Phi_i}{\sqrt{n} \Delta x}, \quad 1 \leq i \leq n, \quad (20)$$

and substitute ρ^∞ into (18), to obtain

$$\lambda_H(\rho^\infty) = \min_{\xi \in \mathbb{R}_+^{n-1}} \left\{ \frac{1}{\Delta x^2} \xi^T A \xi \quad : \quad \xi^T \xi = 1 \right\},$$

where

$$A = \begin{pmatrix} 2 & -1 & & & \\ -1 & 2 & -1 & & \\ & \ddots & \ddots & \ddots & \\ & & -1 & 2 & -1 \\ & & & -1 & 2 \end{pmatrix} \in \mathbb{R}^{(n-1) \times (n-1)}.$$

It is simple to observe that A is positive definite and that¹⁵

$$\begin{aligned} \lambda_H(\rho^\infty) &= \frac{1}{\Delta x^2} \times (\text{the smallest eigenvalue of } A) = \frac{1}{\frac{(b-a)^2}{(n-1)^2}} [2 - 2 \cos(\frac{\pi}{n-1})] \\ &= \frac{\pi^2}{(b-a)^2} + o(1). \end{aligned}$$

Next, we analyze the convergence rate for the cycle graph C_n . Again we assume the relation (19) and let ξ as in (20). Since C_n has one more edge than L_n , we let $\eta \in \mathbb{R}$:

$$\eta := \frac{\Phi_1 - \Phi_n}{\sqrt{n} \Delta x} = \sum_{i=1}^{n-1} \xi_i.$$

Substituting ρ^∞ into (18), we have

$$\begin{aligned} \lambda_H(\rho^\infty) &= \min_{(\xi, \eta) \in \mathbb{R}_+^n} \left\{ \frac{1}{\Delta x^2} [\xi^T A \xi + 2\xi_1 \eta + 2\xi_{n-1} \eta + 2\eta^2] : \right. \\ &\quad \left. \xi^T \xi + \eta^2 = 1, \quad \eta = \sum_{i=1}^{n-1} \xi_i \right\}. \end{aligned} \tag{21}$$

The following transformations reduce (21) to a simpler eigenvalue problem. Let

$$(\xi, \eta)^T = P \xi^T, \quad \text{where} \quad P = \begin{pmatrix} I \\ \mathbf{1}^T \end{pmatrix} \in \mathbb{R}^{n \times (n-1)}$$

with the identity matrix $I \in \mathbb{R}^{(n-1) \times (n-1)}$ and $\mathbf{1} \in \mathbb{R}^{n-1}$ being the vector of all 1's. Then, (21) becomes

$$\lambda_H(\rho^\infty) = \min_{\xi \in \mathbb{R}_+^{n-1}} \left\{ \frac{1}{\Delta x^2} (P \xi)^T B (P \xi) : (P \xi)^T (P \xi) = 1 \right\}, \tag{22}$$

where

$$B = \begin{pmatrix} A & b^T \\ b & 2 \end{pmatrix} \in \mathbb{R}^{n \times n} \quad \text{with} \quad b^T \in \mathbb{R}^{n-1}, \quad b = (1, 0, \dots, 0, 1),$$

and A is as above.

Below, we compute (22). First, we give explicit formulas for the eigenvalues and eigenvectors of B .

¹⁵Here the eigenvector of A corresponding to the smallest eigenvalue satisfies the assumption (19).

Lemma 14. *Let $n \geq 3$. For each $k = 0, 1, \dots, n-1$, the eigenvalues of B are*

$$\lambda_k = 2 - 2 \cos\left(\frac{2k\pi}{n}\right).$$

For $k = 0, 1, \dots, n-1$, the associated eigenvectors in un-normalized form are:

$$v_k = (v_k(j))_{j=1}^n, \quad w_k = (w_k(j))_{j=1}^n,$$

where, for $j = 1, \dots, n-1$,

$$v_k(j) = \sin\left(\frac{2\pi kj}{n}\right), \quad w_k(j) = \cos\left(\frac{2\pi kj}{n}\right);$$

and when $j = n$,

$$v_k(n) = -\sin\left(\frac{2\pi kn}{n}\right), \quad w_k(n) = -\cos\left(\frac{2\pi kn}{n}\right).$$

Proof. The proof is by direct computation. We just show the details for the case of $j = 1$. We have

$$\begin{aligned} (Bv_k)(1) &= 2v_k(1) - v_k(2) + v_k(n) \\ &= 2\sin\left(\frac{2\pi k}{n}\right) - \sin\left(\frac{2 \cdot 2\pi k}{n}\right) - 0 \quad \text{By double angle formula} \\ &= (1 - 2\cos\frac{2k\pi}{n})v_k(1). \end{aligned}$$

And

$$\begin{aligned} (Bw_k)(1) &= 2w_k(1) - w_k(2) + w_k(n) \\ &= 2\cos\left(\frac{2\pi k}{n}\right) - \cos\left(\frac{2 \cdot 2\pi k}{n}\right) + 1 \quad \text{By double angle formula} \\ &= (1 - 2\cos\frac{2k\pi}{n})w_k(1). \end{aligned}$$

□

Note that in Lemma 14, many eigenvalues are repeated. As a consequence, obviously there are only two eigenvectors associated to each repeated eigenvalues, and not four; the repeating eigenvalues, in fact, have identical pairs v_k, w_k , up to sign. However, the eigenvalue equal to 0 is simple, with associated eigenvector $w_0 = (1, \dots, 1, -1)^T$. Moreover, aside from this 0 eigenvalue, all other eigenvalues are positive.

Now, observe that $P^T w = 0$, and therefore the matrix $V = [w, P]$ is invertible and

$$BV = V \begin{bmatrix} 0 & 0 \\ 0 & C \end{bmatrix},$$

where $C \in \mathbb{R}^{n-1, n-1}$. Further, notice that $P^T P$ is positive definite and thus it has a unique positive definite square root $(P^T P)^{1/2}$. Thus, $\xi^T P^T B P \xi$, subject to $(P \xi)^T P \xi = 1$, can be rewritten as

$$\xi^T P^T B P \xi = \xi^T P^T P C \xi = \xi^T (P^T P)^{1/2} (P^T P)^{1/2} C (P^T P)^{-1/2} (P^T P)^{1/2} \xi$$

and thus, with $x = (P^T P)^{1/2} \xi$, we end up with the problem

$$\min_{x: x^T x = 1} x^T \left[(P^T P)^{1/2} C (P^T P)^{-1/2} \right] x.$$

Finally, we notice that the matrix $[(P^T P)^{1/2} C (P^T P)^{-1/2}]$ is symmetric, and it is obviously similar to C , so that indeed

$$\begin{aligned} \min_{x: x^T x = 1} x^T \left[(P^T P)^{1/2} C (P^T P)^{-1/2} \right] x = \\ \min_{\xi \in \mathbb{R}_+^{n-1}} \{ (P\xi)^T B (P\xi) : (P\xi)^T (P\xi) = 1 \} = \text{The second smallest eigenvalue of } B. \end{aligned} \quad (23)$$

Putting it all together, (22) gives

$$\begin{aligned} \lambda_{\mathcal{H}}(\rho^\infty) &= \frac{1}{\Delta x^2} (\text{the second smallest eigenvalue of } B) \\ &= \frac{1}{\frac{(b-a)^2}{(n-1)^2}} [2 - 2 \cos\left(\frac{2\pi}{n}\right)] = \frac{4\pi^2}{(b-a)^2} + o(1), \end{aligned}$$

and the proof of Theorem 13 is completed. \square

From Theorem 13, we see that the graph structure plays a key role in the convergence rate. In particular, the cyclic graph structure in the above model gives a decay rate 4 times that of the lattice case. Although the above model is for the linear entropy, Theorem 7 indicates more than that. The convergence rate $\lambda_{\mathcal{F}}(\rho)$ holds for the entire free energy, which can contain the interaction energy $\frac{1}{2} \sum_{i=1}^n \sum_{j=1}^n w_{ij} \rho_i \rho_j$. The question of “how the graph structure and interaction energy affect the convergence rate” deserves further study.

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